# Models with radiative neutrino masses and minimal dark matter

Guo-Zhu Ning (宁国柱)

北京大学高能物理研究中心

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## Outline

#### 1. Introduction and Motivation

2. Construction of Models

3. Conclusion

## 1. Introduction and Motivation



## 2013 NOBEL PRIZE IN PHYSICS

## **Francois Englert**

## Peter W. Higgs

PPT from Q.H. Cao's talk

"for the theoretical discovery of a mechanism that contributes to our understanding of the origin of mass of subatomic particles, and which recently was confirmed through the discovery of the predicted fundamental particle, by the ATLAS and CMS experiments at CERN's Large Hadron Collider"

SM is not the whole story

neutrino mass, DM

### Oscillation experiments





SNO

Kam LAND

other atmospheric, solar, accelerator, and reactor neutrino experiments



Neutrino have tiny mass and flavor violation

## **Evidence for Dark Matter** at Different Astrophysics Scales

**Spiral Galaxies** Scale: ~10<sup>21</sup> m Rotation curves remain flat far beyond the edge of the visible disk.

$$\begin{array}{lll} v(R) &=& \sqrt{GM(R)/R} \\ v(R) &\approx & const \end{array} \right\} \Rightarrow \left\{ \begin{array}{ll} M(R) &\propto & R \\ \rho(R) &\propto & R^{-2} \end{array} \right.$$

#### Galaxy Clusters

Scale: ~1022 m

14

0.8

0.6

0.4

0.2

0.2

Ω,

- Orbital velocities of galaxies (Zwicky's discovery in 1933)
- X-ray gas
- Gravitational lensing

#### The Dark Universe - Scale: ~10<sup>26</sup> m

- CMB: Ω<sub>tot</sub>=1.0
- CMB, BBN: Ω<sub>b</sub>=0.045
- Galaxy clusters: Ω<sub>m</sub>=0.27
- Supernovae la:  $\Omega_m$ ,  $\Omega_\Lambda$
- Structure formation: cold DM



## Neutrino in the SM and beyond

In SM, neutrino is massless :

(1) Gauge symmetry;

(2) Renormalization;

(3) Minimal particle spectrum; No RH neutrino, only one Higgs doublet

Possible corresponding ways to extend:

(1) Extension of gauge group: eg. SUSY, Grand Unification.

(2) Effective operator: eg. Weinberg operator.

(3) Extension of particle spectrum:

(a) Canonical seesaw; (b) Quantum effects.



### DM: a model view

#### (1) Relic density;

WIMP: a particle with weak interaction naturally give relic desity with mass at electroweak scale

#### (2) Stable Or decay lifetimes larger than the age of the Universe;

Symmetry, higher dimensional operator.....

#### Stabilization mechanism in SM

Photon: U(1)QED; Neutrino: Lorentz invariance;

#### (3) Others;

(a) electroneutrality;

(b) Y=0 from DM direct detection experiment;

(c) Non-baryonic;

electron: lightest charged particle; proton: baryon number conservation



#### One stone two birds



(1) Additional Z2 is imposed to keep the stability of DM;

Minimal DM may work for this problems

(2) A relatively small number of paramters due to (accident) symmetry;

## Minimal DM

(M. Cirelli et al, NPB 753(2006)178, NPB 787(2007)152, NJP 11(2009)105005)

#### Basic idea:

Add to SM one extra multiplets with minimal spin, isospin and hypercharge quantum numbers, and satisfy :

(1) The lightest component is automatically stable; | Higher isospin

(2) The DM is still allowed by DM searches;

hyperchargeless

#### Main conclusions:

(1)  $n \ge 5$ , for fermions,  $n \ge 7$  for scalars;

(2) Small number of parameters:

(a) For fermions, gauge type interactions and one new parameter( the tree-level mass);

(b) For scalars, quartic couplings with the Higgs fields besides the same thing like the case in fermions.

#### **Motivation**

Radiative neutrino mass + Minimal DM

(1) Radiative neutrino mass :

With the help of higher isospin multiplets we do a systematic investigation, from dimensional-5 Weinberg operator, on the neutrino mass generated at one loop.

(2) The Stability of DM:

Analysis of the symmetry lead by the assignments of the quantum number of new fields. eg. Z2 symmetry

## 2. Construction of models

### One-loop realization of Weinberg operator



Analogous Type I/III



Only considerate fermions and scalars running in the loop, not include color particles and vector bosons.

Analogous Type II



(1) The divergence can be absorbed by the tree level counter-terms coming from the regular tree level Type II (Type II-like) seesaw;
 (2) To link DM, the fields running in the loop should have neutral component.

## Lesson from minimal DM model

Fermion multiplets

Mass term, gauge interactions

$$T \ge \frac{3}{2}, n = 2T + 1 \ge 4, U(1)$$
 Sym

Neutrality and Y=0

(1) Symmetry; 
$$T \ge 2, n \ge 5, U(1)$$
$$Q = T3 + Y, T3 \subset [-T, T]$$

Scalar multiplets

Mass term, gauge interactions, scalar potential

 $T \ge 3, n \ge 7, Z2$ 

$$T \ge \frac{5}{2}, n = 2T + 1, Y = 0$$
, Z2 Symmetry;

Neutrality and Y=0

Generalization on Scalar multiplets

(part of results: K. Earl et al. arxiv:1311.3656)

(1)  $Y = 0, T \ge 3, T$  odd integer, Z2 (2)  $T \ge Y \ge 1, T \ge 2, Y, T$  integer, U(1)

(3)  $T \ge Y \ge 1, T \ge 5/2, Y, T$  half integer, U(1)

Notice: Y can be nonzero in generalization

### Constraints on multilplets

#### Perturbativity of SU(2)L gauge coupling

#### Model dependent

- $n \leq 5$  , for Majorana fermions
- $n \leq 6$  , for complex scalar

 $n \le 8$ , for real scalars (Planck scale)  $n \le 7$ , for real scalars (TeV scale)

#### Perturbative unitarity

 $n \leq 8$  , for complex scalar

New fields in our analysis

 $n \leq 9$ , for real scalars (TeV scale)

 $Q = T3 + Y, T3 \subset [-T, T]$ 

Scalars (1) Y = 0, T = 3, for real scalar; (2)  $Y \neq 0, \frac{1}{2}, T = \frac{5}{2}, \frac{7}{2}$ , for complex scalar; (3)  $Y \neq 0, T = 2,3,4$ , for complex scalar.

fermions 
$$T = \frac{3}{2}, 2, \frac{5}{2}$$
, for real scalar;

Y can set any values only if there exists neutral component in multiplets

#### An elaborate example

Lepton number violation in the mass term

$$N: T = 2, Y = 0;$$
  
 $\eta: T = \frac{5}{2}, \frac{3}{2}, Y = \frac{1}{2}; \ \widetilde{H}\eta\eta\widetilde{\eta}$ 

Lepton number violation in the Yukawa coupling

$$N: T = 2, Y = 2, 1;$$
  

$$\eta_1: T = \frac{5}{2}, Y = \frac{5}{2}, \frac{3}{2};$$
  

$$\eta_2: T = \frac{5}{2}, Y = \frac{3}{2}, \frac{1}{2}; \widetilde{H}\eta_2\eta_2\widetilde{\eta_1}$$
  

$$N: T = \frac{5}{2}, Y = \frac{5}{2}, \frac{3}{2}, \frac{1}{2};$$
  

$$\eta_1: T = 3, Y = 3, 2, 1;$$
  

$$\eta_2: T = 3, Y = 2, 1, 0; \widetilde{H}\eta_2\eta_2\widetilde{\eta_1}$$



**Careful:** When more than one new scalar fields are introduced, we must make sure that the new interactions brought by them don't break the pre-existing symmetry.

## **Relevant phenomenology**

(1) Flavor physics will be relevant if new fermion is not to heavy and coupling can be large;

(2) New scalar may have influence in Higgs physics ;

(3) Due to exotic charge, new fields will have rich collider phenomenology if their masses are not heavy ;

## 3. Conclusion

#### Conclusion

(1) A systematic investigation has been done for the models with radiative neutrino masses and minimal dark matter;

(2) radiative neutrino masses and minimal dark matter can be achieved in a simple model except hyperchargeless;

(3) These models will have rich phenomenology if the new fields is not too heavy and its couplings with SM particles not too weak ;

## Thanks for your attention !!!