

Search for FCNC with top quark and Higgs at ATLAS experiment



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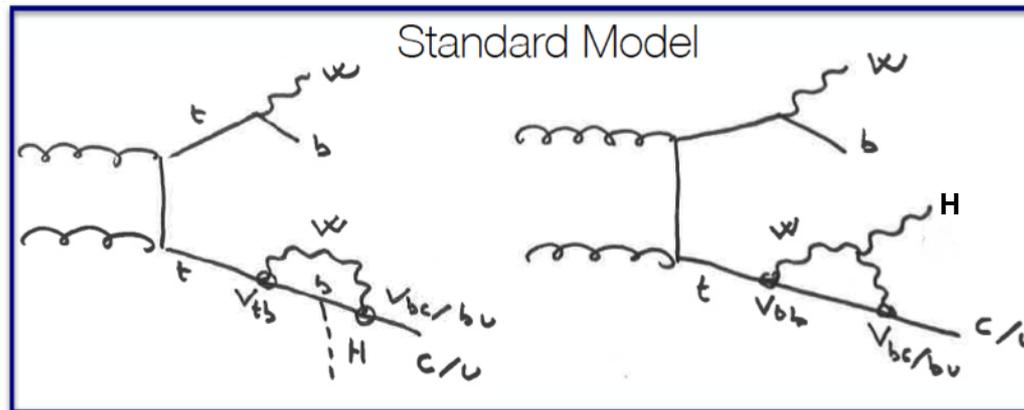


粒子物理前沿讨论会

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Introduction

- Flavor changing neutral current (FCNC) in the decay of $t \rightarrow Hq$ ($q=u,c$) is highly suppressed in the SM
 - The GIM can render partial cancelation among the loop amplitudes where the Higgs is irradiated by a W



$$V_{tb} V_{cb}^* + V_{ts} V_{cs}^* + V_{td} V_{cd}^*$$

- BR can be enhanced by BSM:

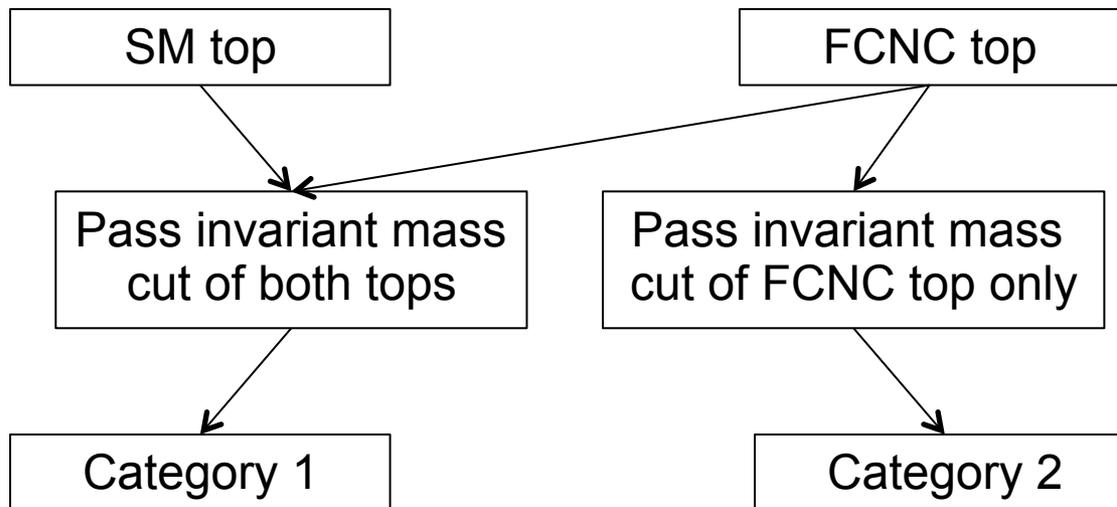
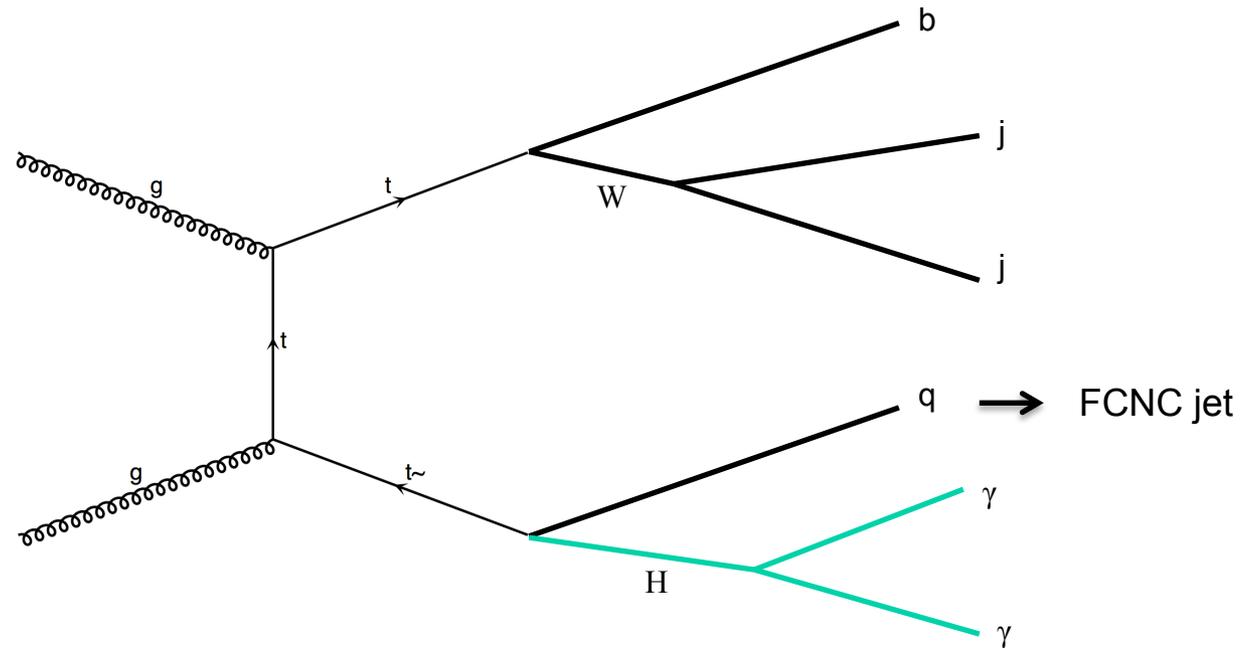
	SM	QS	2HDM	FC 2HDM	MSSM	RPV SUSY
BR($t \rightarrow Hu$)	2×10^{-17}	$\sim 10^{-5}$	$\sim 10^{-5}$	-	$\sim 10^{-5}$	$\sim 10^{-6}$
BR($t \rightarrow Hc$)	3×10^{-15}	$\sim 10^{-5}$	$\sim 10^{-3}$	$\sim 10^{-5}$	$\sim 10^{-5}$	$\sim 10^{-6}$

[PRD 67 (2003) 035003, PRD 55 (1997) 3156, PRD 75 (2007) 075021, PRD 58 (1998) 055001, PRD 75 (2007) 015002, et al.]

FCNC $t \rightarrow qH(\gamma\gamma)$ search

[JHEP 10 (2017) 129]

Events with 2 photons ,
 $p_T > 40, 30$ GeV,
 $100 < m_{\gamma\gamma} < 160$ GeV

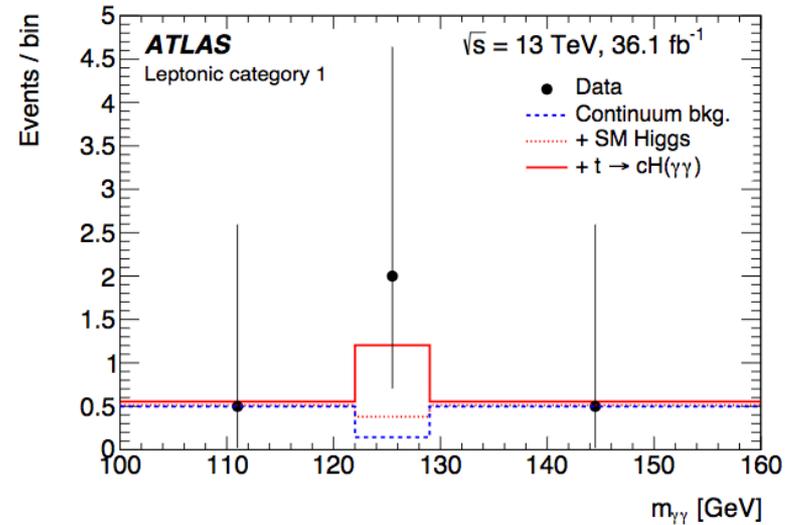
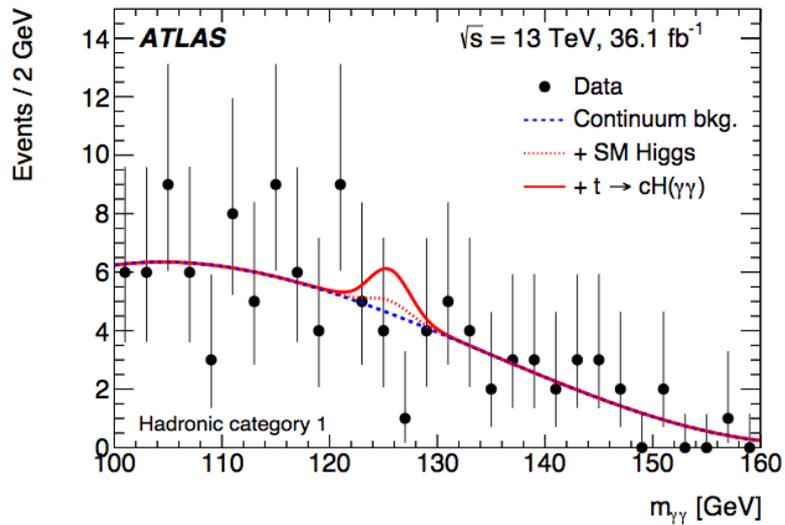
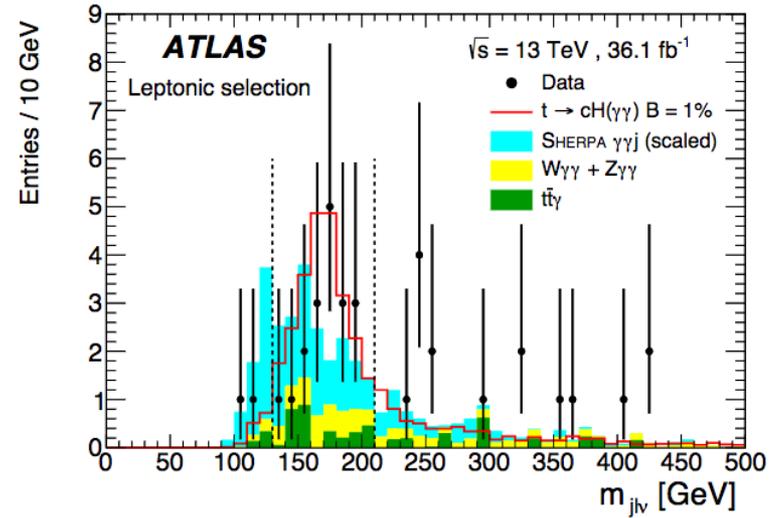
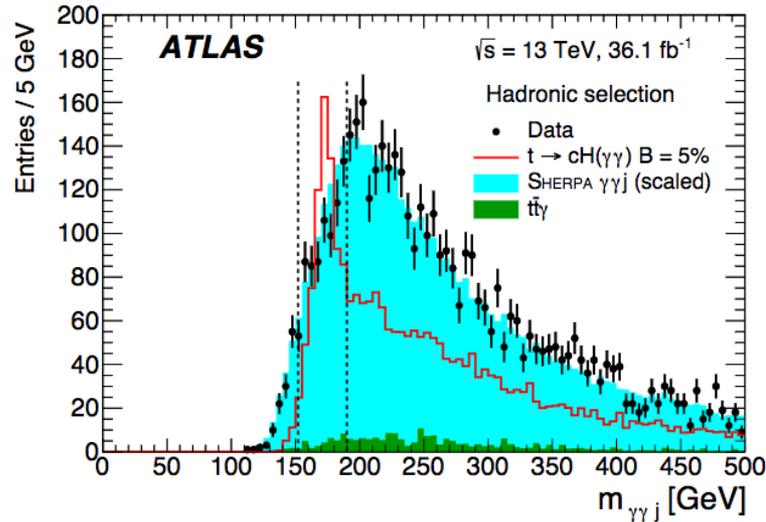


Events are further divided into leptonic and hadronic modes depending on the other top decay

4/2 jets for hadronic/leptonic channels

FCNC $t \rightarrow qH(\gamma\gamma)$ search

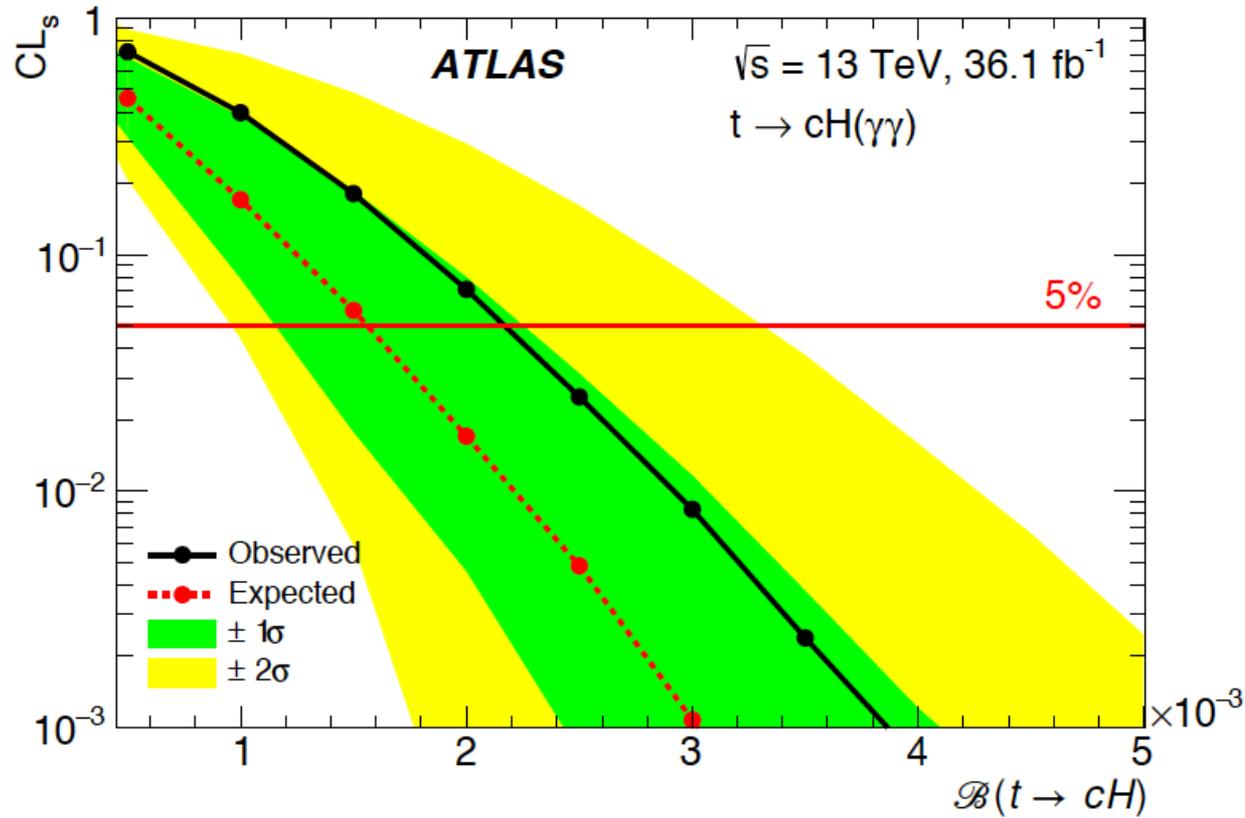
[JHEP 10 (2017) 129]



Fit the $m_{\gamma\gamma}$ spectra to extract the sensitivity

FCNC $t \rightarrow qH(\gamma\gamma)$ search

[JHEP 10 (2017) 129]



Obs. (exp.) BR($t \rightarrow uH$) (%)	Obs. (exp.) BR($t \rightarrow cH$) (%)
0.24 (0.17)	0.22 (0.16)

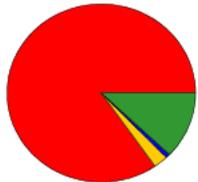
FCNC $t \rightarrow qH$ with multilepton (ML) search

- Same production, but search for $H \rightarrow \tau\tau, WW^*, ZZ^*$ with 2 or 3 leptons (e/μ) in the final state
 - Same-sign 2 leptons, ≥ 4 jets, 1 or 2 b-jets
 - 3 leptons, ≥ 2 jets, ≥ 1 b-jet

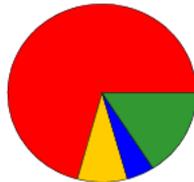
ATLAS Simulation
 $\sqrt{s} = 13$ TeV



2 ℓ SS



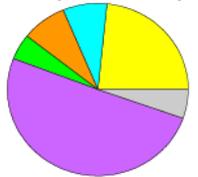
3 ℓ



ATLAS Simulation
 $\sqrt{s} = 13$ TeV, 36.1 fb⁻¹



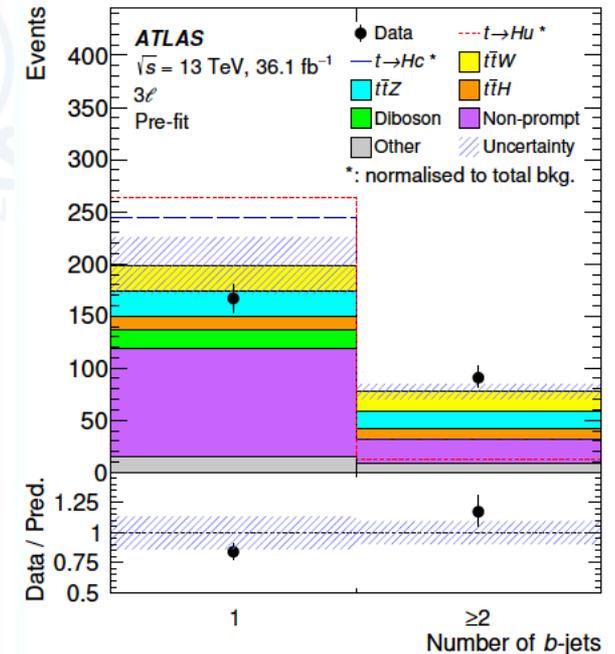
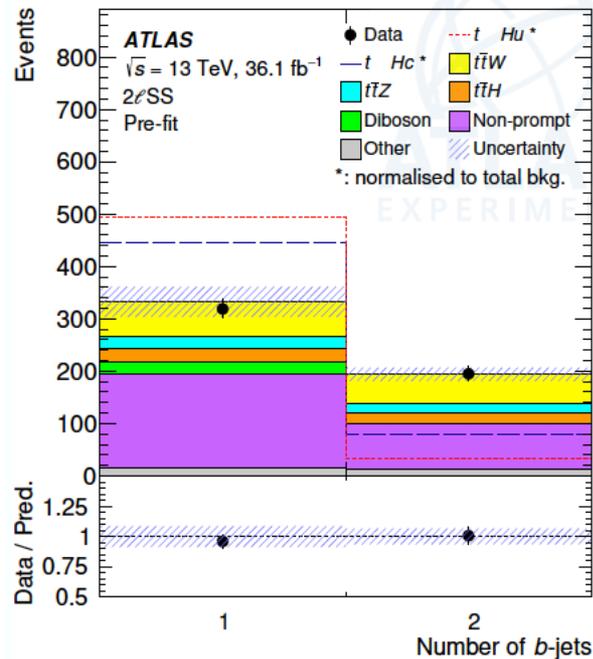
2 ℓ SS (526 events)



3 ℓ (276 events)



[PRD 98 (2018) 032002]

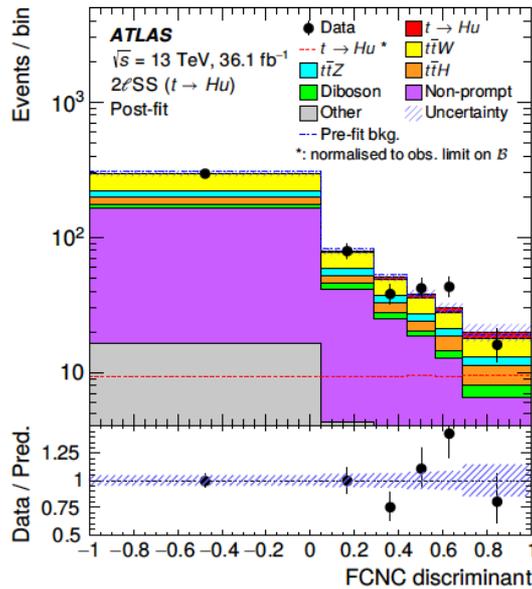


FCNC $t \rightarrow qH$ with multilepton (ML) search

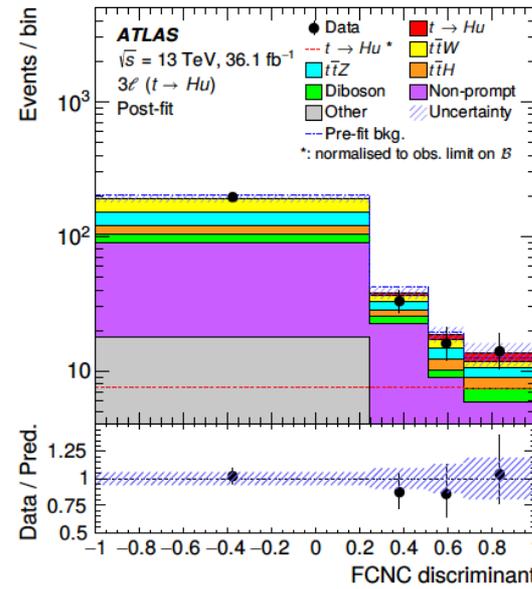
- Re-use ttH analysis, update new background estimation and new BDT training

[PRD 98 (2018) 032002]

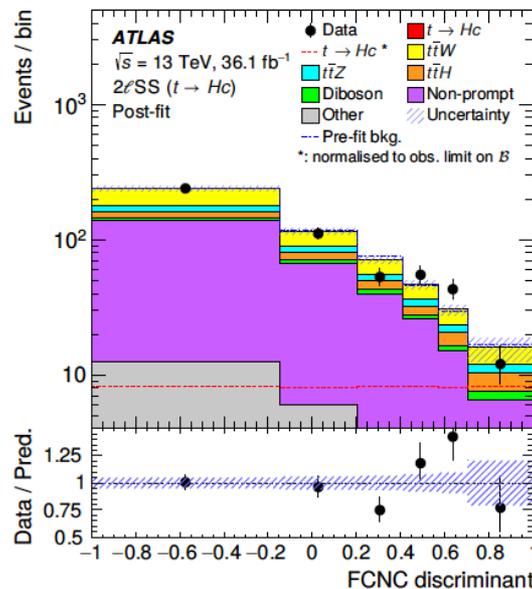
2l SS
 $t \rightarrow uH$



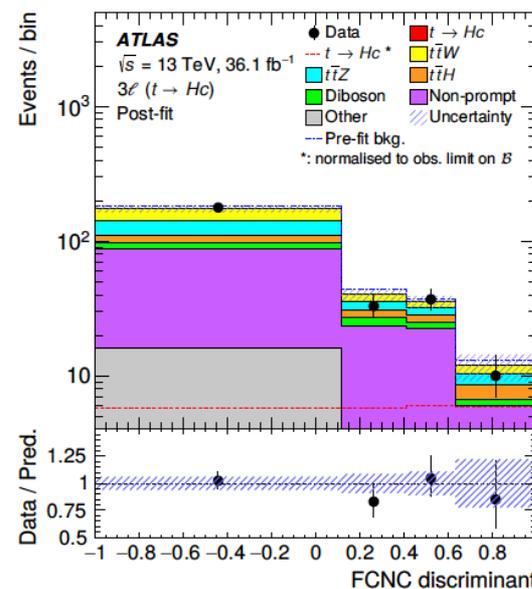
3l
 $t \rightarrow uH$



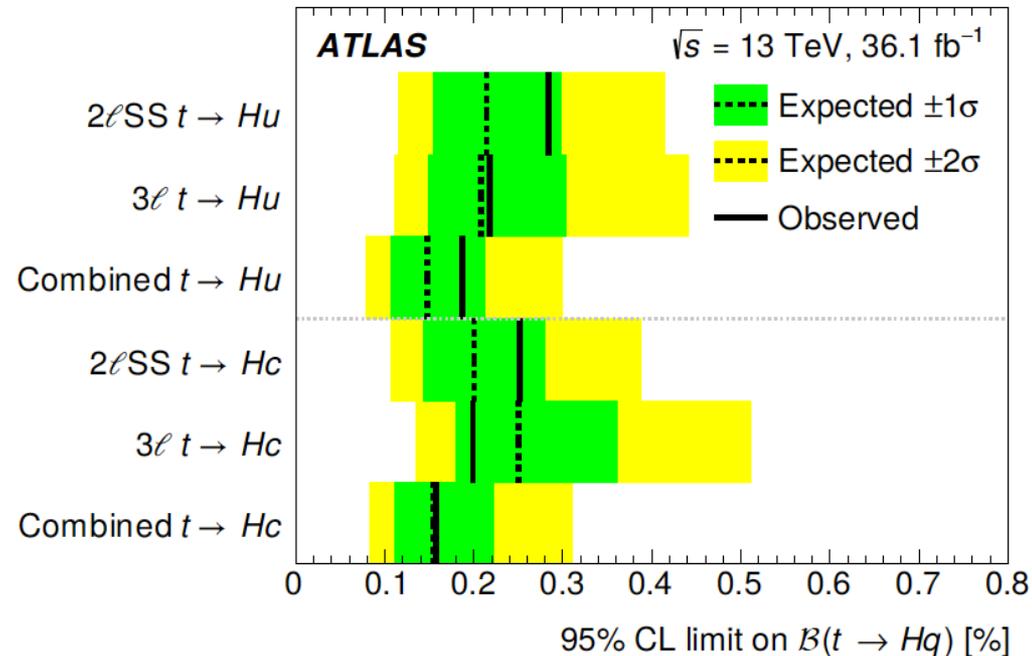
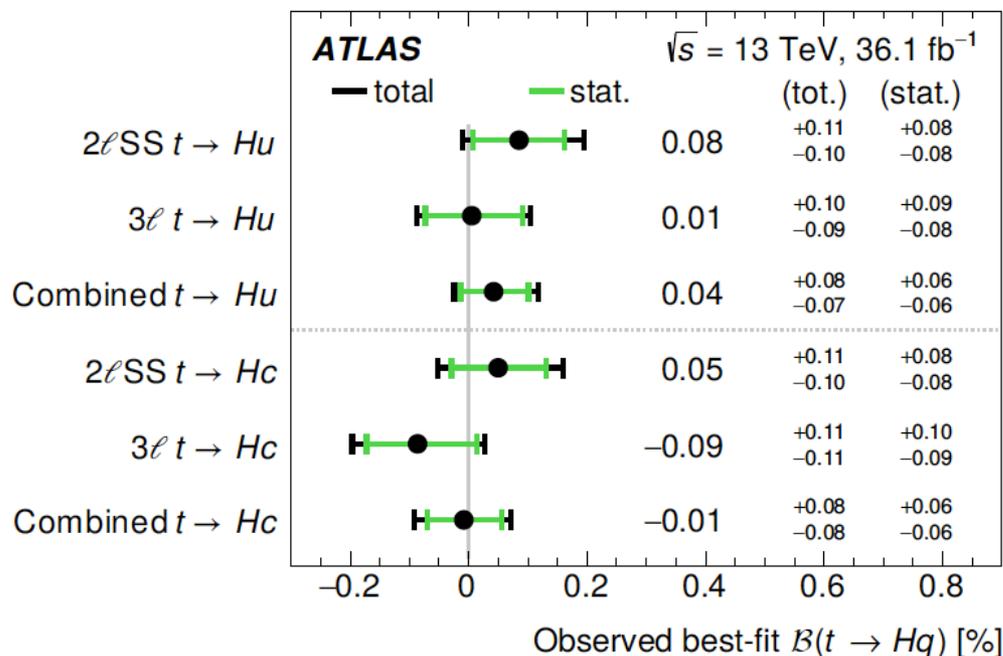
2l SS
 $t \rightarrow cH$



3l
 $t \rightarrow cH$



FCNC $t \rightarrow qH$ with multilepton (ML) search



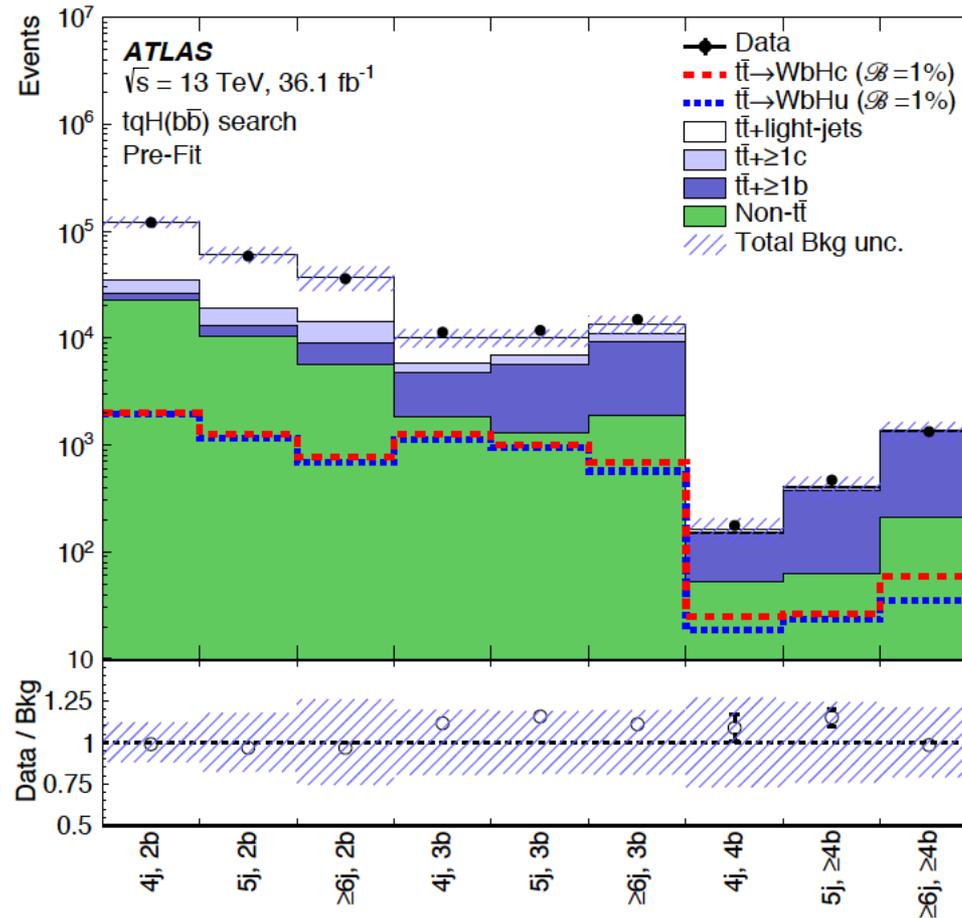
[PRD 98 (2018) 032002]

Obs. (exp.) BR($t \rightarrow uH$) (%)	Obs. (exp.) BR($t \rightarrow cH$) (%)
0.16 (0.15)	0.19 (0.15)

FCNC $t \rightarrow qH(bb)$ search

[arXiv:1812.11568]

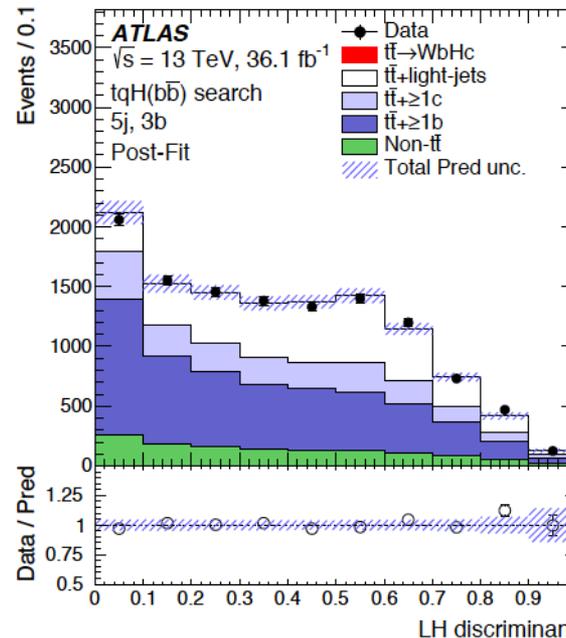
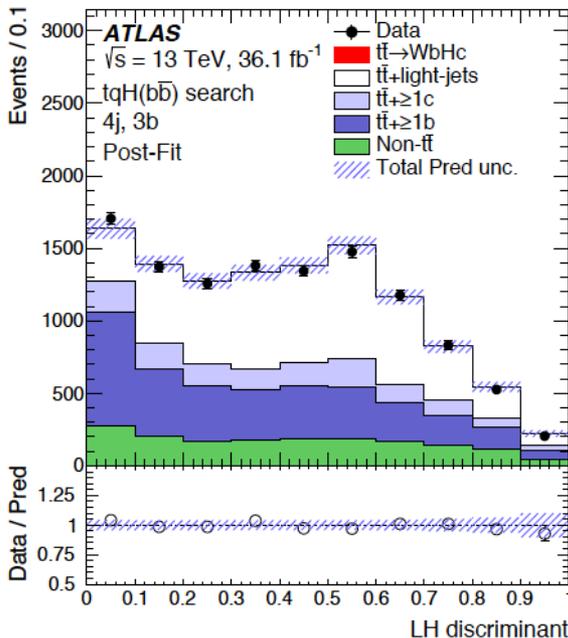
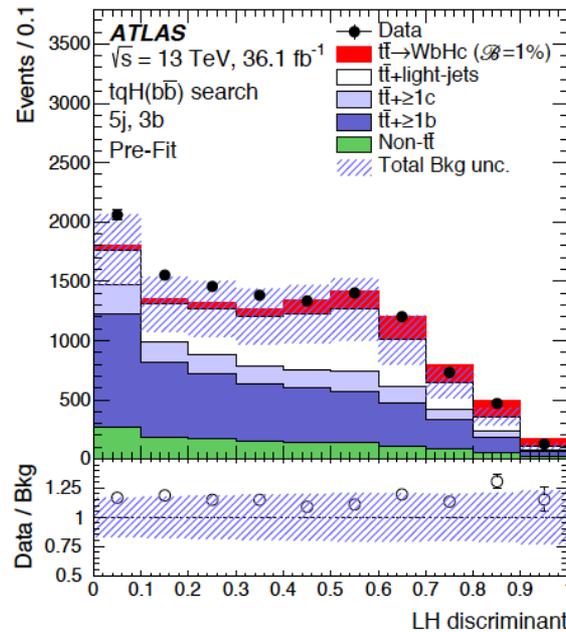
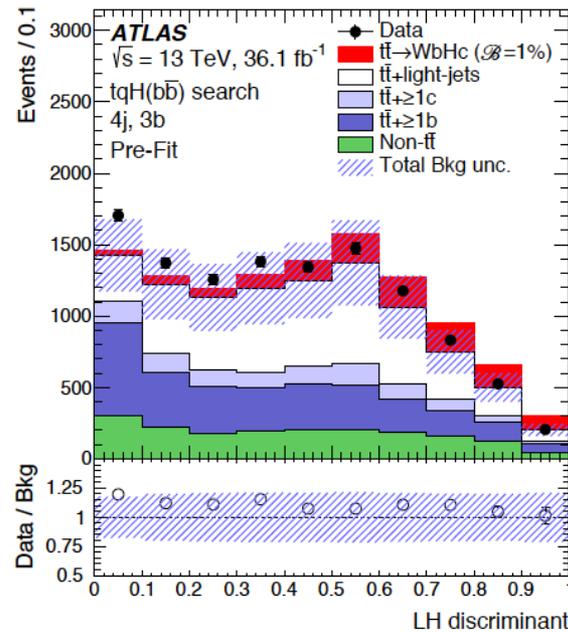
- Require exactly 1 lepton, ≥ 4 jets, ≥ 2 b-jet
- 9 analysis regions with different combinations of n jets + m b-jets



- Most sensitive regions are 5j,3b and 4j,3b

FCNC $t \rightarrow qH(bb)$ search

[arXiv:1812.11568]

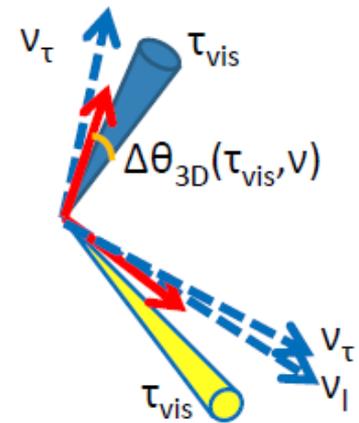
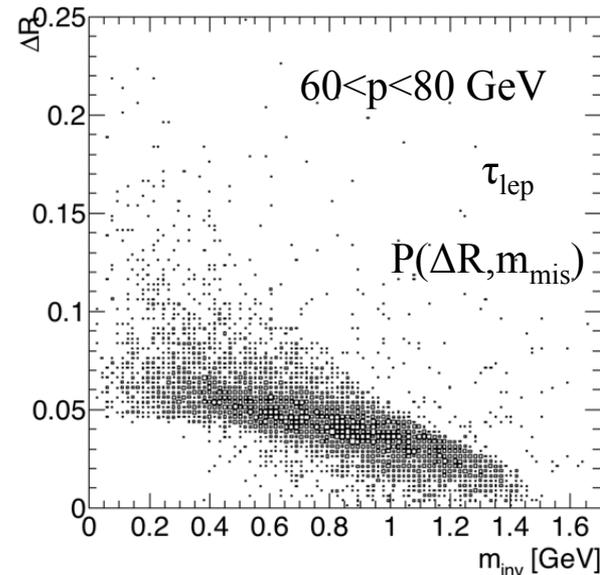
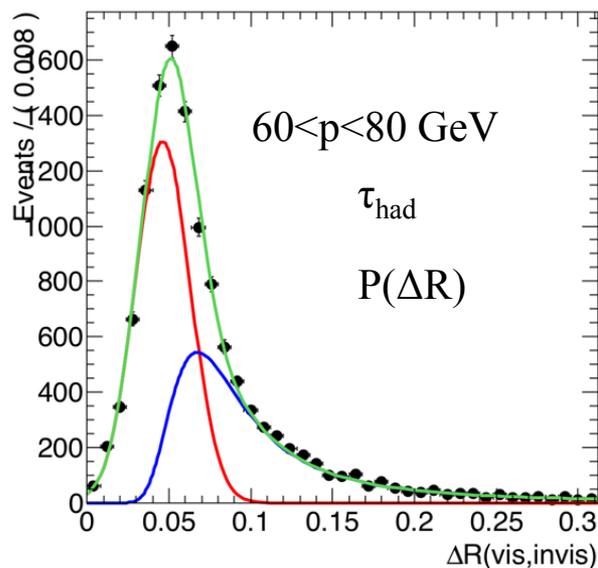


- Use Likelihood (LH) discriminant based on two-body and three-body invariant masses that correspond to the expected resonances under signal and background hypotheses per event

$$L(\mathbf{x}) = \frac{P^{\text{sig}}(\mathbf{x})}{P^{\text{sig}}(\mathbf{x}) + P^{\text{bkg}}(\mathbf{x})}$$

FCNC $t \rightarrow qH(\tau\tau)$ search

- Due to boosted taus, the neutrinos from tau decay tend to be aligned with the visible tau decay products. Assuming exact collinearity, the di-tau mass can be uniquely reconstructed (“collinear mass”)
- However, ATLAS used a mass reconstruction method, which builds a likelihood that incorporates tau decay kinematics – MMC, which improves the di-tau mass reconstruction



Hadronic and leptonic taus are parameterized (based on MC info.) differently, depending on if the total invisible 4-vector has a mass

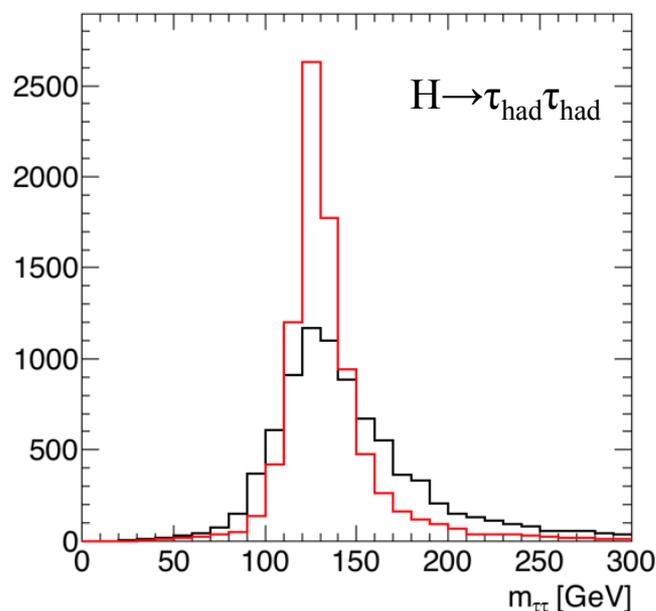
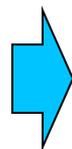
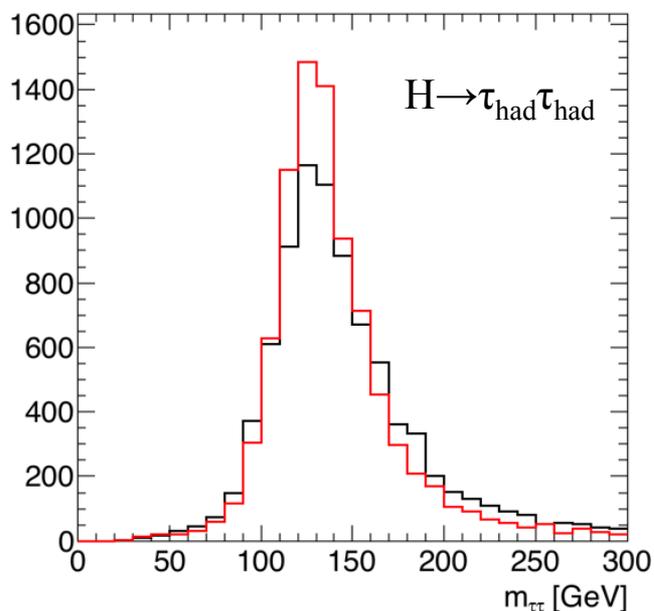
H $\rightarrow\tau\tau$ Mass

[PRD 93 (2016) 113010]

- The invisible 4-vectors from tau decays are obtained per event by minimizing a combined χ^2 :

$$\chi^2 = \underbrace{-2 \ln P_1 - 2 \ln P_2 + \left(\frac{m_{\tau_1}^{\text{fit}} - 1.78}{\sigma_\tau} \right)^2 + \left(\frac{m_{\tau_2}^{\text{fit}} - 1.78}{\sigma_\tau} \right)^2}_{\text{Tau Kinematics Constraint}} + \underbrace{\left(\frac{m_H^{\text{fit}} - 125}{\sigma_H} \right)^2}_{\text{Higgs Mass}} + \underbrace{\left(\frac{E_{x,\text{mis}}^{\text{fit}} - E_{x,\text{mis}}}{\sigma_{\text{mis}}} \right)^2 + \left(\frac{E_{y,\text{mis}}^{\text{fit}} - E_{y,\text{mis}}}{\sigma_{\text{mis}}} \right)^2}_{\text{MET Constraint}}$$

The H $\rightarrow\tau\tau$ mass resolution improves with different constraints:

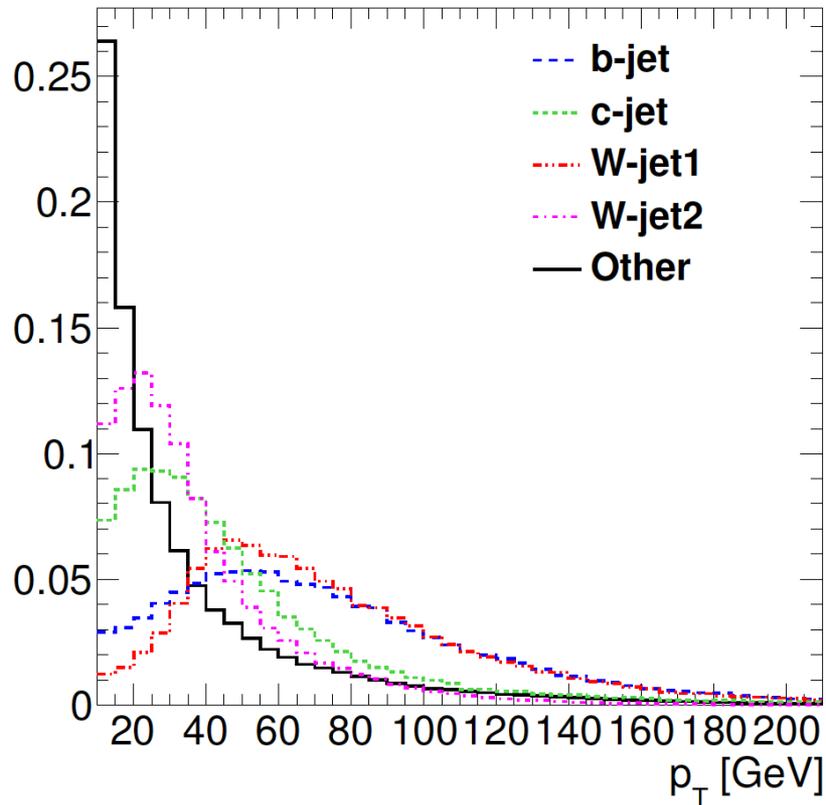


Black: MET constraint only (collinear mass)

Red (left): MET + tau kinematics constraint

Red (right): MET + tau kinematics constraint + Nominal Higgs mass constraint

Signal categorization ($t \rightarrow qH, H \rightarrow \tau\tau$) [PRD 93 (2016) 113010]



Assigning jets to the signal topology ($t\bar{t} \rightarrow WbHc \rightarrow qqbt\tau c$) can not always be done correctly: some jets' p_T is below object acceptance, multiple extra jets from QCD radiation:

- Radiation jets are often softer than the ones from top decay
- The FCNC jet from $t \rightarrow Hc$ decay, and the subleading jet from W boson decay, have non-negligible chances of missing the object selection, with the former leading to **unmatched** part of the signal

Consequently, divide the events into 3 jets and ≥ 4 jets categories, and each is further divided into $\tau_{lep}\tau_{had}$ and $\tau_{had}\tau_{had}$ channels

Jet assignment ($t \rightarrow qH, H \rightarrow \tau\tau$)

[PRD 93 (2016) 113010]

In the so-called 3-jet category:

- For the 3-jet events (at least one of which should be b-tagged), if they can form a top, and satisfy

$$\left(\frac{m_{j_1 j_2} - 80.4}{20} \right)^2 + \left(\frac{m_{j_1 j_2 b} - 172.5}{25} \right)^2 < 5,$$

Then this event is discarded (no c-jet in the event)

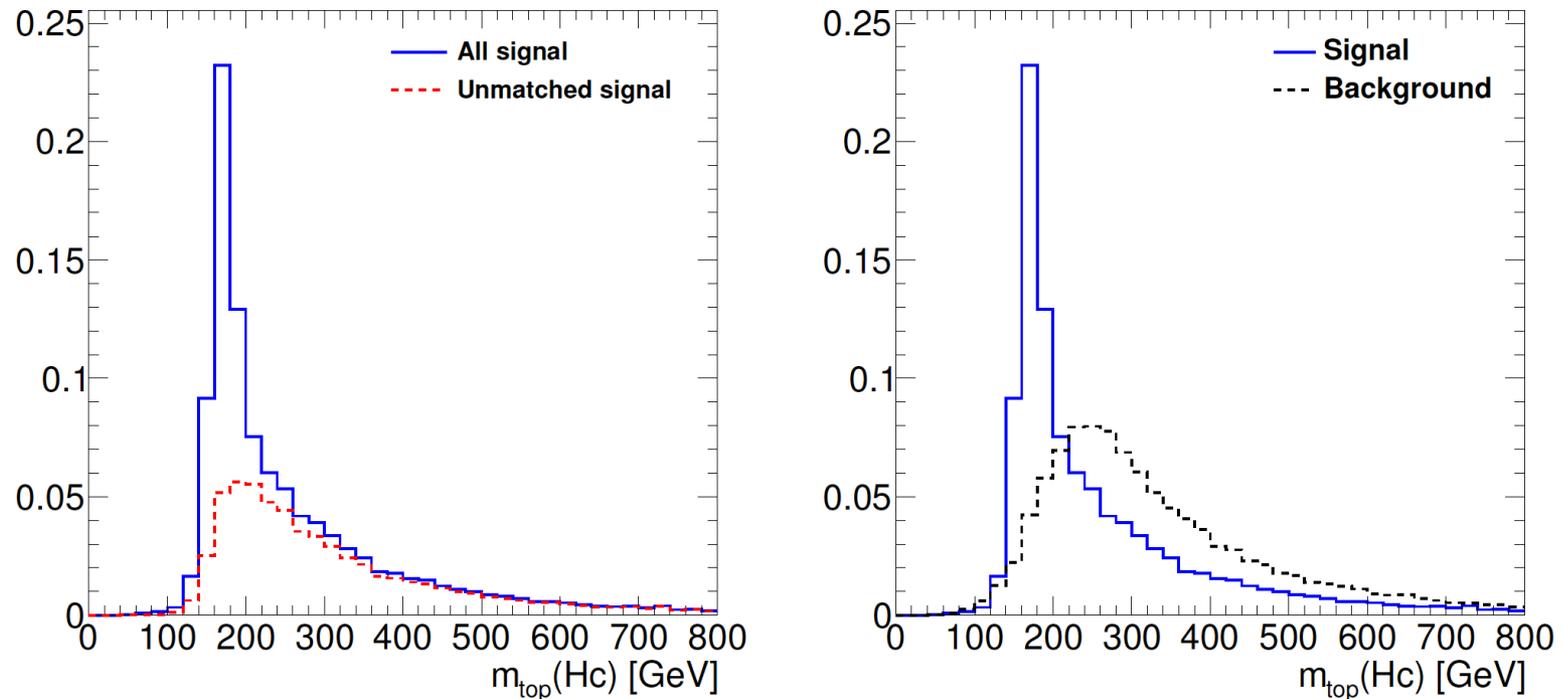
- Otherwise, the c-jet is defined as the configuration that gives the least ΔR sum of (think of the boosted topology):

$$\Delta R(j_c, H) + \Delta R(j_{W1}, b)$$

In the 4-jet bin (≥ 4 jets inclusively), only the leading 4 jets are considered, out of which at least one should be b-tagged. The c-jet is also found by a minimum ΔR sum:

$$\Delta R(j_c, H) + \Delta R(j_{W1}, b) + \Delta R(j_{W2}, b) + \Delta R(j_{W1}, j_{W2})$$

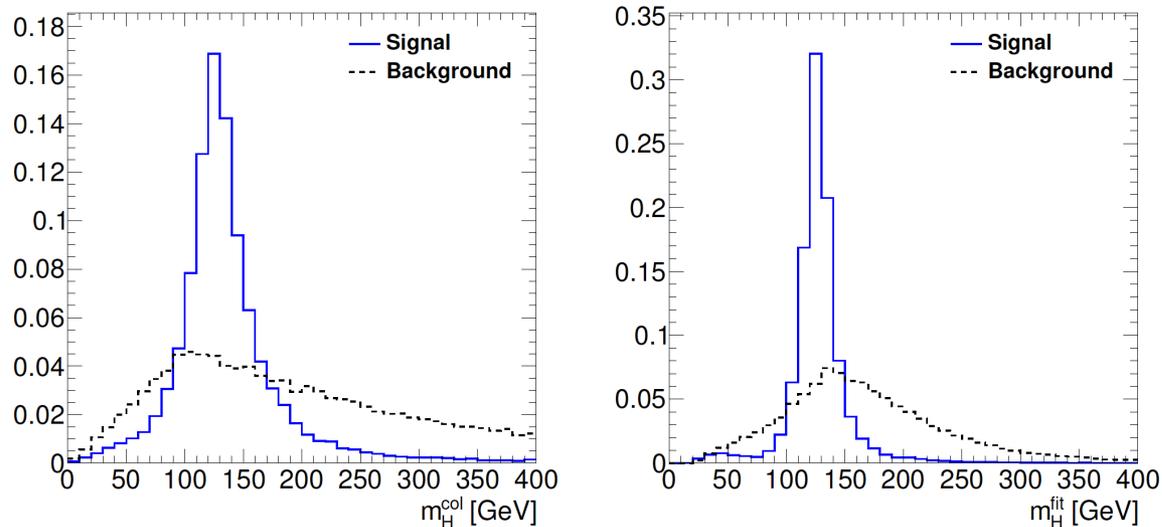
After the jet assignment, the top mass from the $t \rightarrow Hc$ decay can be reconstructed:



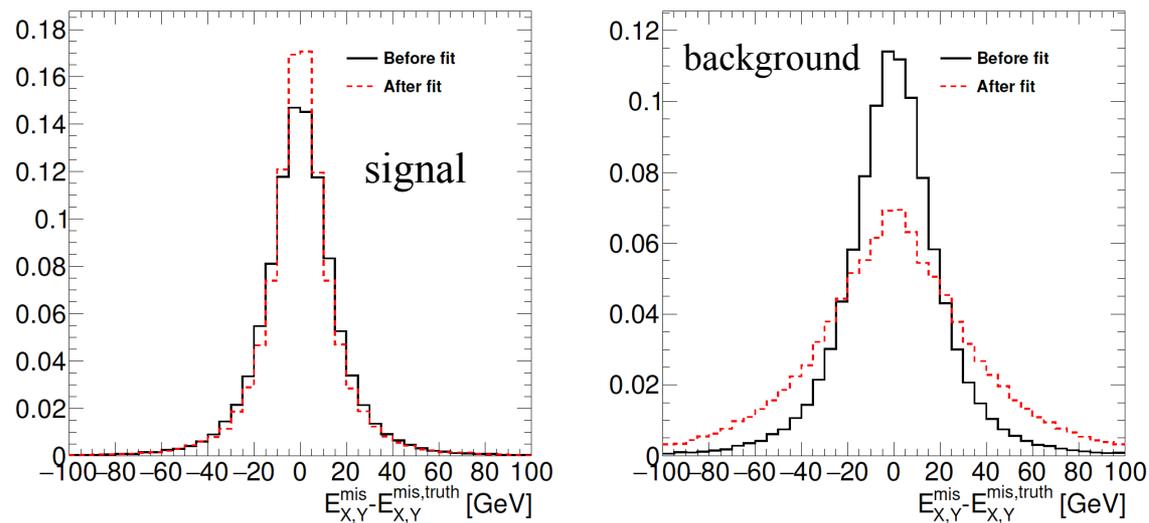
- Analogous to the $H \rightarrow \gamma\gamma$ channel, a top mass peak appears for events with truth-matched FCNC jet, but we also have a fraction of the signal events with unmatched FCNC jet
- The background is also not peaking in this mass distribution as shown in the right plot

Higgs mass and MET ($t \rightarrow qH$, $H \rightarrow \tau\tau$) [PRD 93 (2016) 113010]

The di-tau mass in background (dominated by $t\bar{t}$) also changes after the fit:

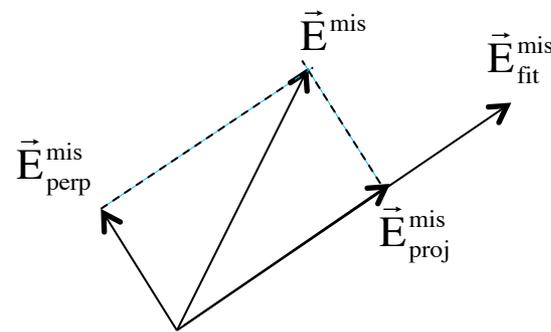
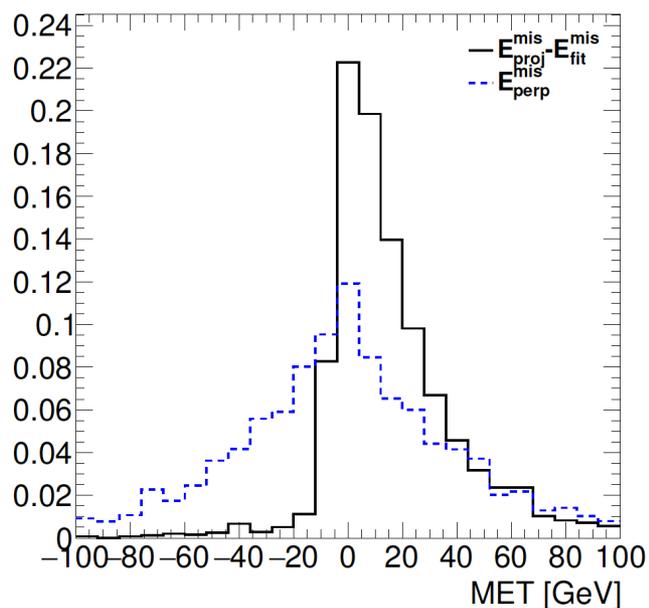
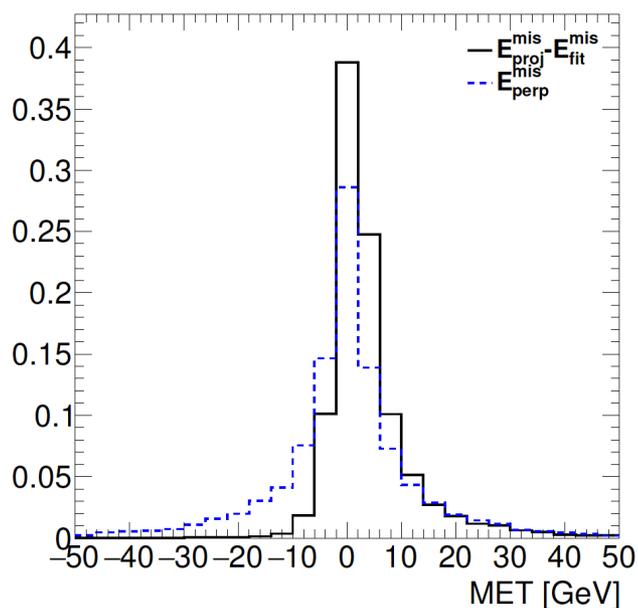


Meanwhile, the MET gets worse (better) in the background (signal):



Higgs mass and MET ($t \rightarrow qH$, $H \rightarrow \tau\tau$) [PRD 93 (2016) 113010]

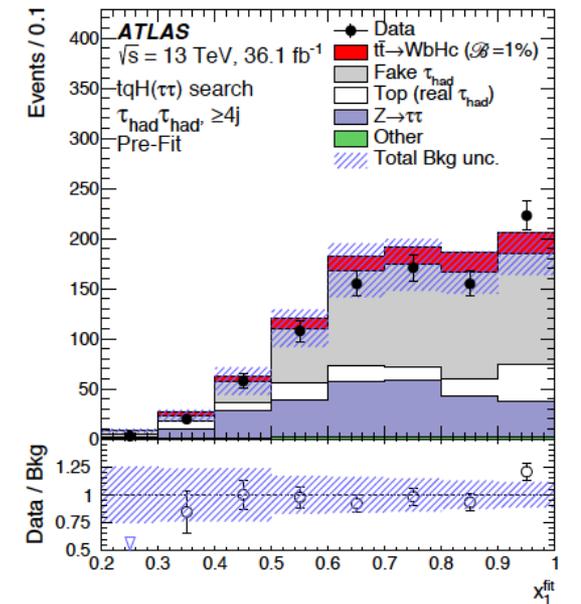
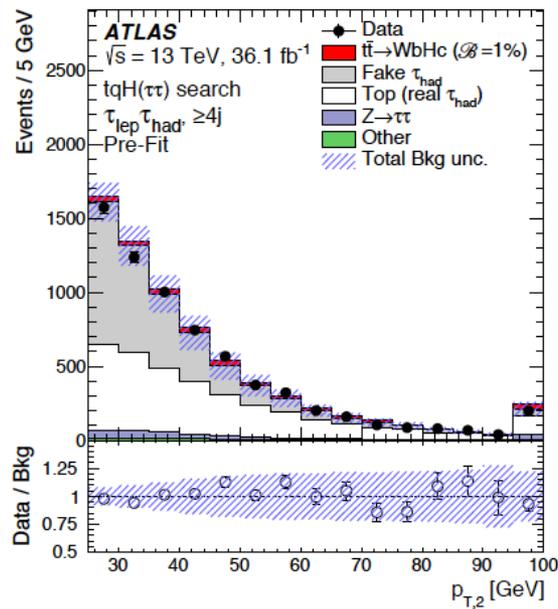
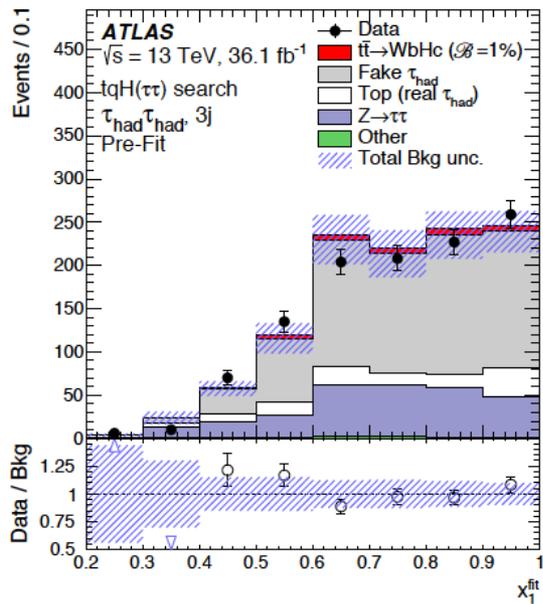
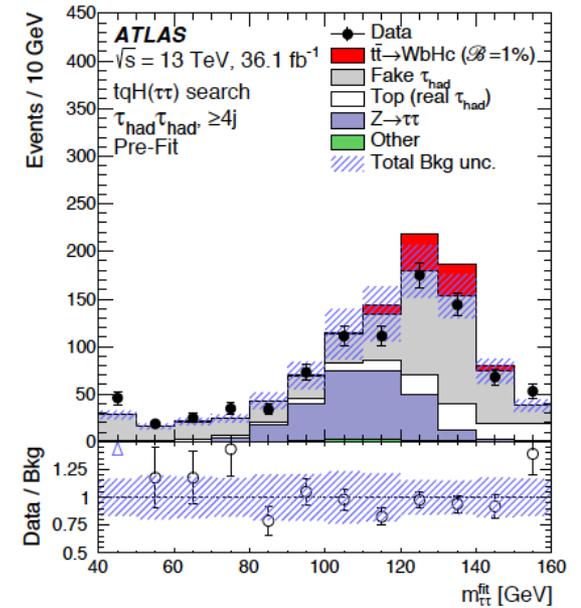
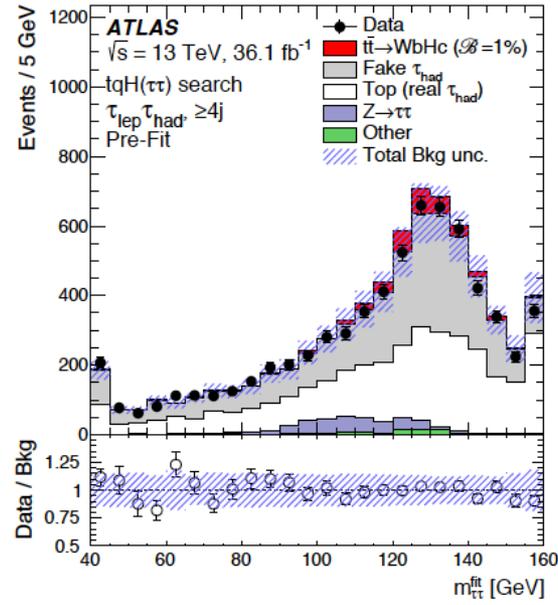
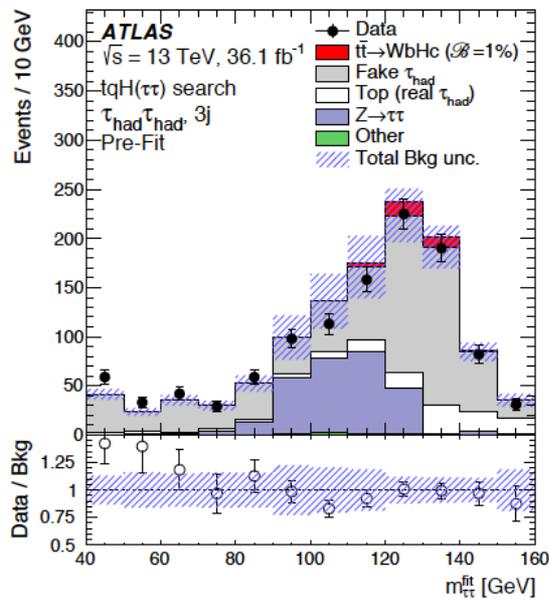
- At reconstruction level, this means that MET varied within its resolution to get closer to the true MET in signal, while it has to vary a lot (a large fake MET is created) to conform to the mass constraint
- The variation in MET vector can be expressed in terms of MET projections:



The MET change can be used as a discriminant between signal and background, with the BDT method

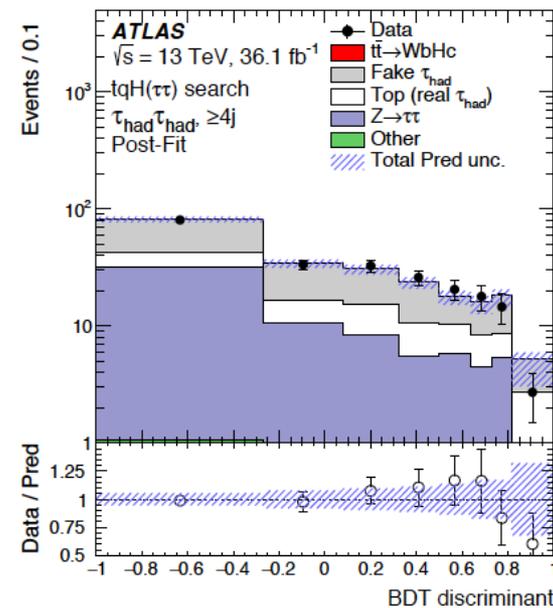
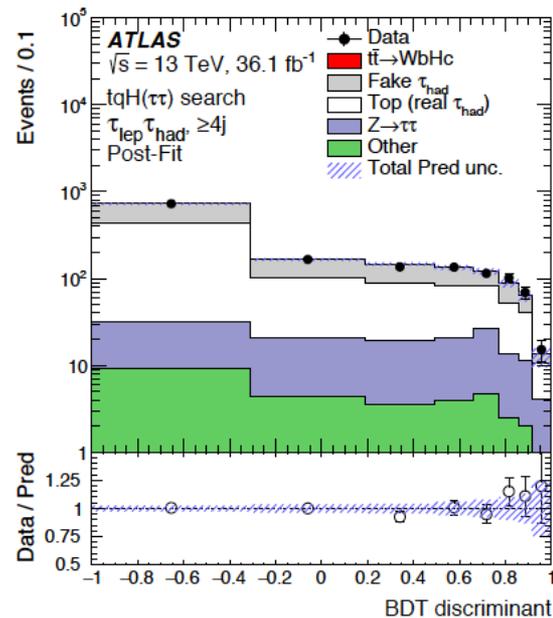
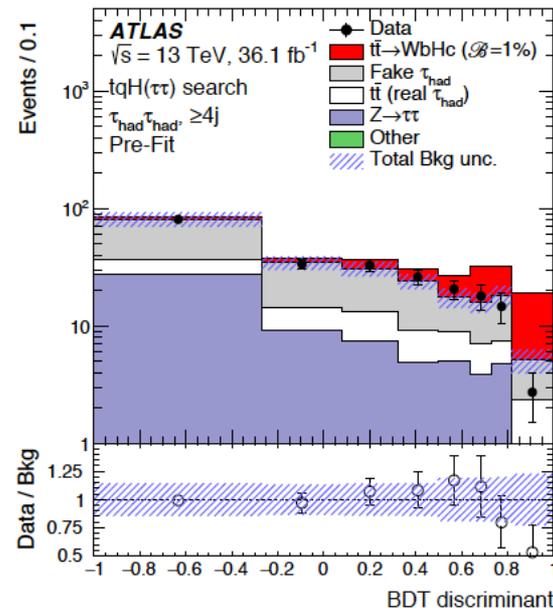
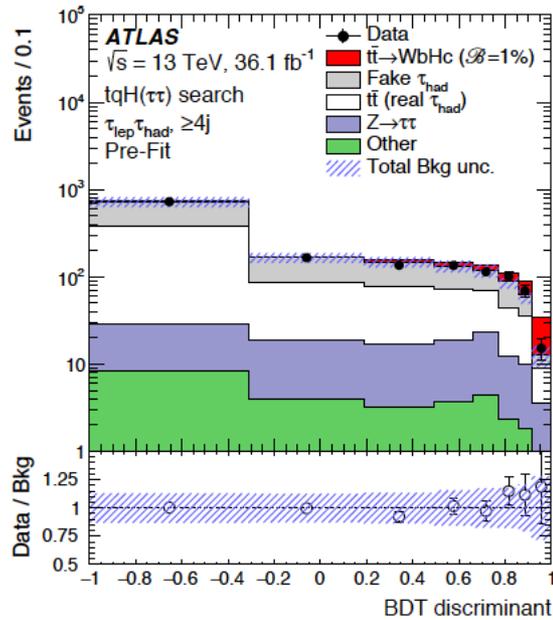
FCNC $t \rightarrow qH(\tau\tau)$ search

[arXiv:1812.11568]



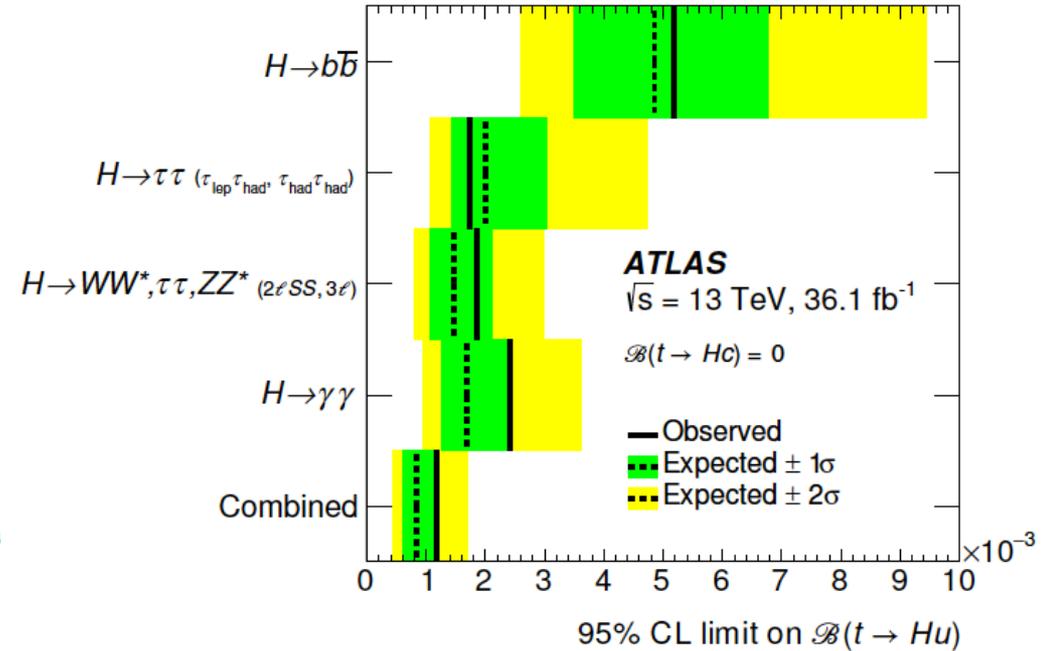
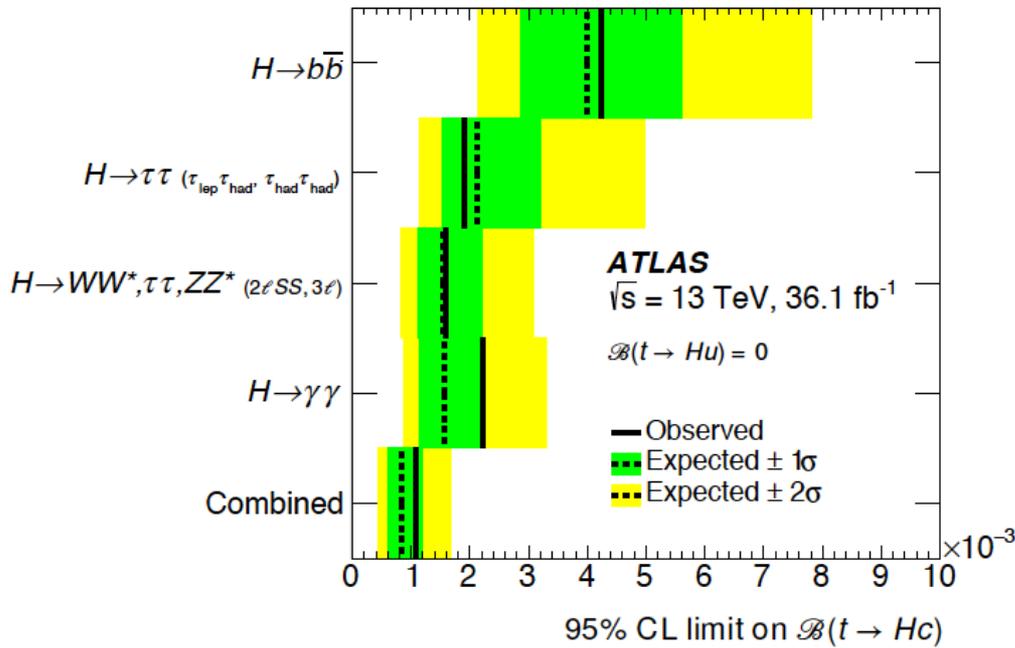
FCNC $t \rightarrow qH(\tau\tau)$ search

[arXiv:1812.11568]



- Require exactly 1 b-jet
- Use data-driven background estimation for jet faking τ_{had}
- Use 9-14 variables for BDT analysis
- Most sensitive regions are 4-jet categories

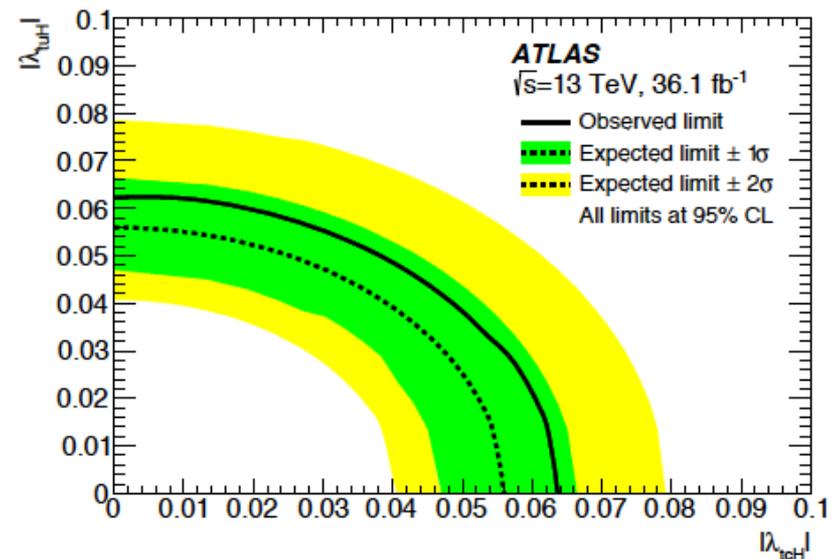
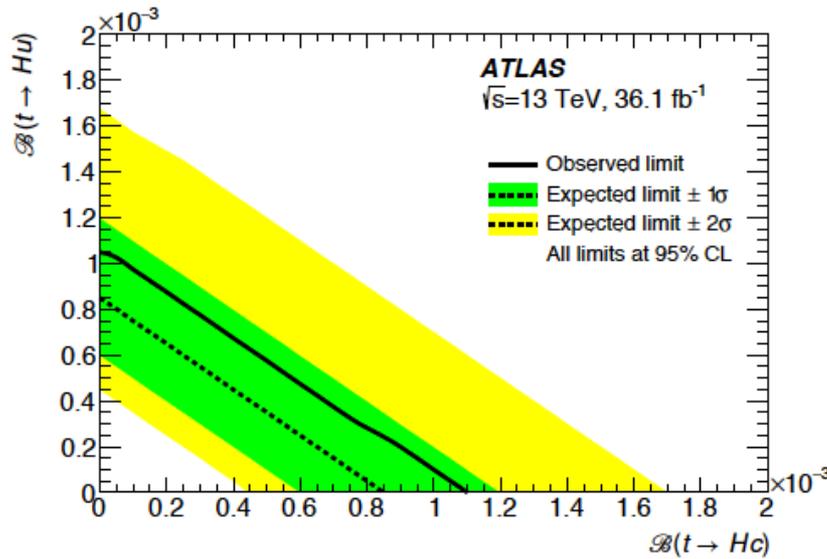
FCNC $t \rightarrow qH(bb, \tau\tau)$ search [arXiv:1812.11568]



	Obs. (exp.) BR($t \rightarrow uH$) (%)	Obs. (exp.) BR($t \rightarrow cH$) (%)
$H \rightarrow b\bar{b}$	0.52 (0.49)	0.42 (0.40)
$H \rightarrow \tau\tau$	0.17 (0.20)	0.19 (0.21)

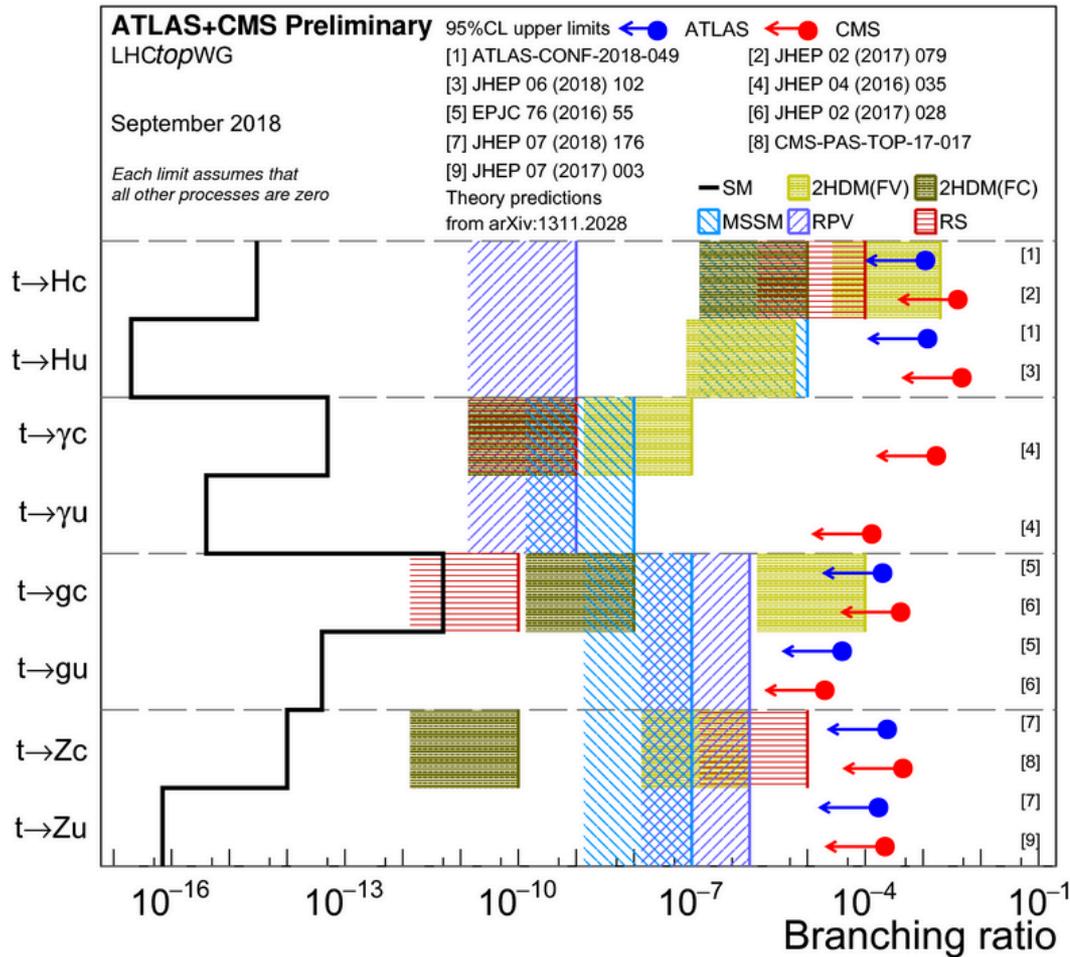
FCNC $t \rightarrow qH$ combination [arXiv:1812.11568]

	95% CL upper limits on $\mathcal{B}(t \rightarrow Hc)$	95% CL upper limits on $\mathcal{B}(t \rightarrow Hu)$
	Observed (Expected)	Observed (Expected)
$H \rightarrow b\bar{b}$	4.2×10^{-3} (4.0×10^{-3})	5.2×10^{-3} (4.9×10^{-3})
$H \rightarrow \tau\tau$ ($\tau_{\text{lep}}\tau_{\text{had}}, \tau_{\text{had}}\tau_{\text{had}}$)	1.9×10^{-3} (2.1×10^{-3})	1.7×10^{-3} (2.0×10^{-3})
$H \rightarrow WW^*, \tau\tau, ZZ^* (2\ell\text{SS}, 3\ell)$ [22]	1.6×10^{-3} (1.5×10^{-3})	1.9×10^{-3} (1.5×10^{-3})
$H \rightarrow \gamma\gamma$ [21]	2.2×10^{-3} (1.6×10^{-3})	2.4×10^{-3} (1.7×10^{-3})
Combination	1.1×10^{-3} (8.3×10^{-4})	1.2×10^{-3} (8.3×10^{-4})

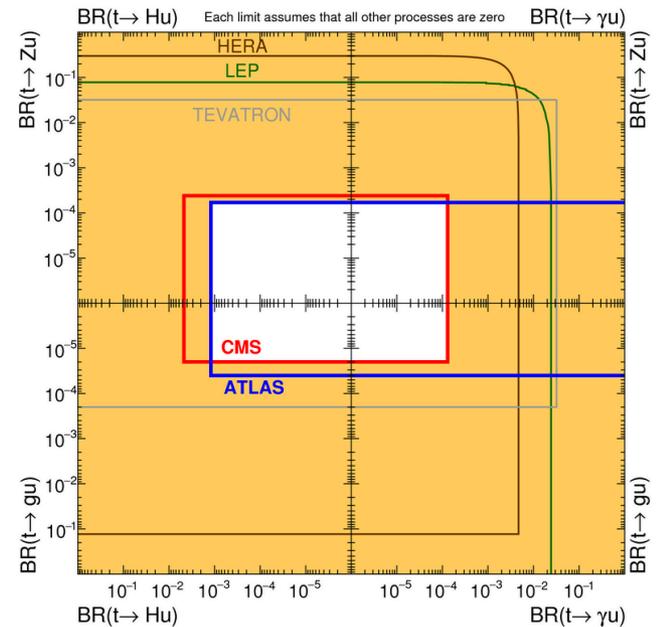
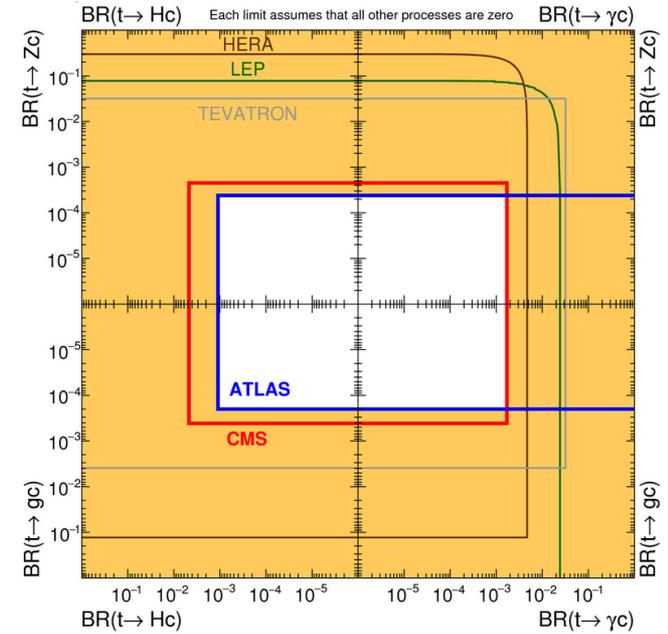


$$\mathcal{L}_{\text{FCNC}} = -\lambda_{tLqR} \bar{t}_L q_R H - \lambda_{qL tR} \bar{q}_L t_R H + h.c. \quad |\lambda_{tqH}| = (1.92 \pm 0.02) \sqrt{\mathcal{B}(t \rightarrow Hq)}$$

LHC top FCNC (summary)



<https://twiki.cern.ch/twiki/bin/view/LHCPhysics/LHCtopWGSummaryPlots>



Summary

First phase of ATLAS program for FCNC top-quark decay into Higgs searches is carried out for all major Higgs decay modes with 36.1 fb^{-1} of run-2 data. No significant excess in $t \rightarrow qH$ process is observed in the $t\bar{t}$ production

The current results give the best limit on $\text{BR}(t \rightarrow qH)$ so far, which also outperforms run-1 limits. Limits entering a region where the most optimal estimation for $\text{BR}(t \rightarrow qH)$ at $\sim 10^{-3}$ as predicted by some models

Search for FCNC continues to provide a very promising area to test BSM physics predictions. Stay tuned for the results based on full Run-2 data, and run-3

Backup Slides

References:

ATLAS: <https://twiki.cern.ch/twiki/bin/view/AtlasPublic/TopPublicResults>